

# Broken Rotational Invariance in AdS/CFT

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# What is AdS/CFT?

- AdS = “Anti de Sitter”
  - A maximally symmetric solution to Einstein’s field equations in vacuum
- CFT = “Conformal Field Theory”
  - Quantum Field Theory that is symmetric under conformal transformations
- AdS/CFT correspondence
  - $d$ -dimensional CFTs have dual descriptions in an  $d+1$ -dimensional AdS space
  - CFT lives on the boundary of the AdS space (holographic)

# Why is AdS/CFT interesting to us?

- Strongly coupled field theories are challenging
- Strongly-coupled field theory = Weak curvature gravitational theory
- Use techniques from GR to study field theories
- AdS/CMT

# AdS/CMT

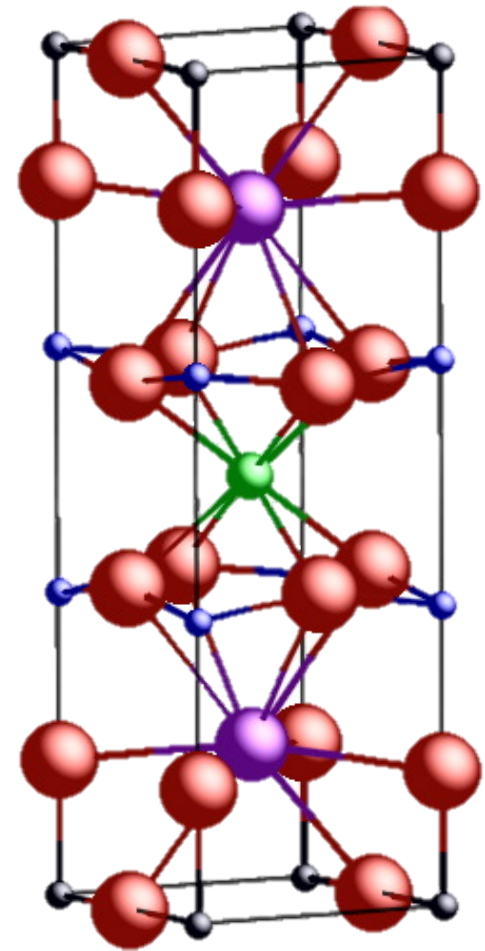
- Are condensed matter systems appropriately described by a conformal field theory?

$$i\hbar \frac{\partial}{\partial t} \Psi(x, t) = \left[ \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x, t) \right] \Psi(x, t)$$

- No
  - Dilation symmetry broken
  - Lorentz symmetry broken

# AdS/CMT

- Can we deform AdS/CFT for condensed matter applications?
  - Lifshitz Geometry
- Is this good enough?
  - We still have rotational symmetry in spatial dimensions



# AdS/CMT

- Spatial rotational invariance not a general property of condensed matter systems
- We need to break this symmetry in our geometry also
- Brings us to our geometry of interest
  - Homogeneous in  $t, x, y$
  - Anisotropic in  $t, x, y$
  - Scaling/Bulk dimension given by  $r$

$$ds^2 = \frac{-dt^2}{r^{2z}} + \frac{dx^2}{r^{2p}} + \frac{dy^2}{r^{2q}} + \frac{dr^2}{r^2}$$

# Energy Conditions

- Used to restrict to physical solutions of Einstein's Field Equations
  - Mass-Energy density is nonnegative
- Field Equations relate matter distribution to geometry, so we can use the Field Equations and the energy conditions to constraint  $\rho$ ,  $p$ ,  $q$
- INSERT CONSTRAINTS HERE

$$T^{\mu\nu} = \begin{bmatrix} \rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{bmatrix}$$

$$\text{NEC: } \rho + p \geq 0$$

$$\text{WEC: } \rho \geq 0, \rho + p \geq 0$$

# Klein-Gordon Equation

- Use Lagrangian Density

$$\mathcal{L} = \sqrt{-g} \left[ \frac{1}{2} g^{\mu\nu} \partial_\mu \partial_\nu \phi - \frac{1}{2} m^2 \phi^2 \right]$$

- Vary with respect to field

$$\partial_\mu \left[ \frac{\partial \mathcal{L}}{\partial_\mu \phi} \right] - \frac{\partial \mathcal{L}}{\partial \phi} = 0$$

- Find equation of motion

$$\frac{-1}{\sqrt{-g}} \partial_\mu \left( g^{\mu\nu} \sqrt{-g} \partial_\nu \phi \right) + m^2 \phi = 0$$



# Schrodinger Potential

- We find a 2<sup>nd</sup> order ODE for our field

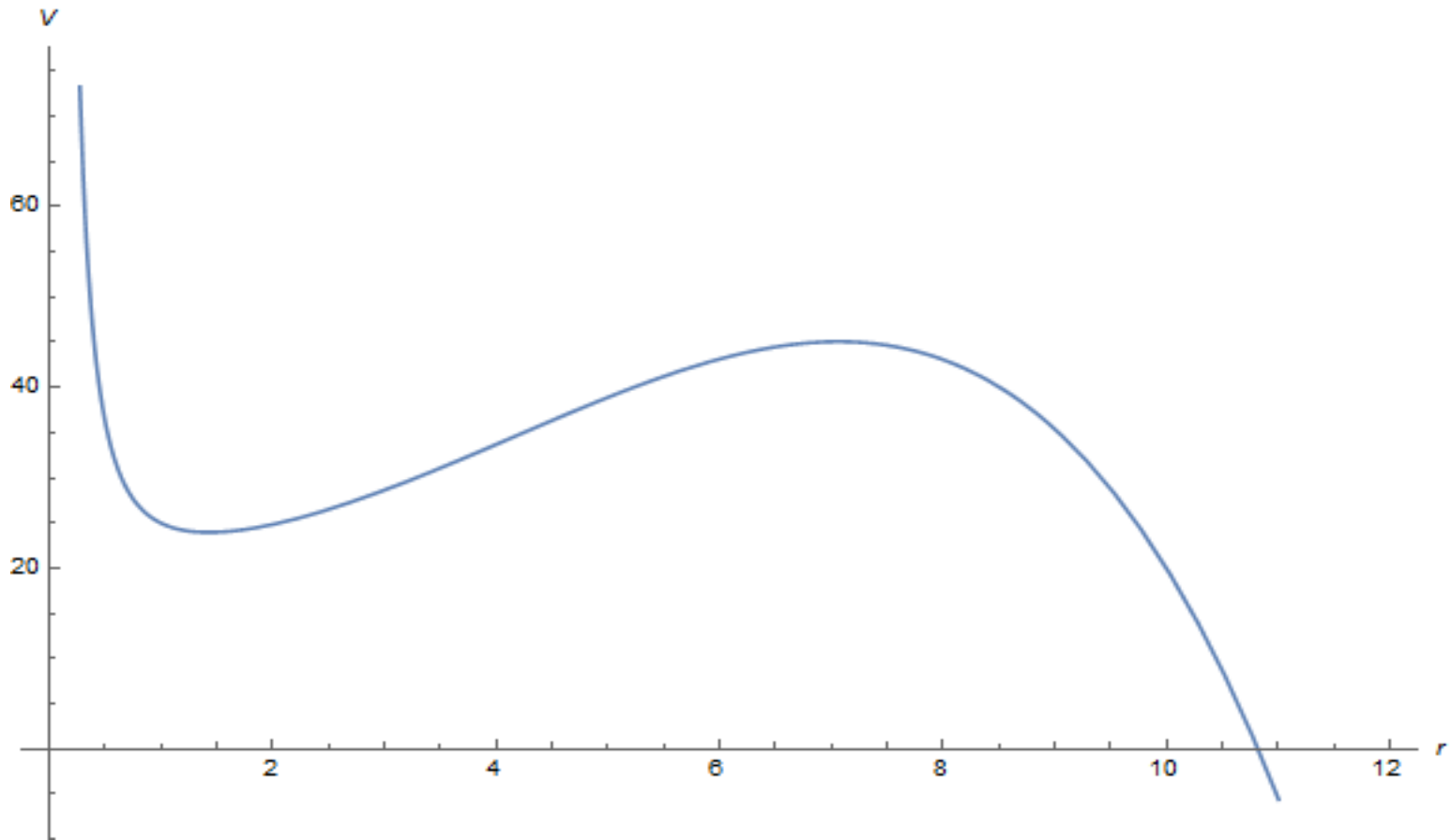
$$-\phi''(r) + \alpha(r)\phi'(r) + \beta(r)\phi(r) = 0$$

- An appropriate field redefinition provides

$$-\phi''(r) + \left( \frac{\nu^2 + k_x^2 r^{2p} + k_y r^{2q} - \omega^2 r^{2z}}{r^2} \right) \phi(r) = 0$$

- This looks like the Time-Independent Schrodinger Equation!
  - Can use techniques from Quantum Mechanics

# Sample Schrodinger Potential



$$V(r) = \frac{\nu^2 + k_x^2 r^{2p} + k_y r^{2q} - \omega^2 r^{2z}}{r^2}$$

# WKB and Green Functions Results

$$\phi_{WKB}(r) = C_{\pm} \frac{1}{|E - V(r)|^{\frac{1}{4}}} e^{\pm i \int_{r_0}^r \sqrt{E - V(r')} dr'}$$

- WKB approximation allows us to understand behavior past the classical turning point
- Prescription from AdS/CFT enables us to find Green Function for field based on behavior at  $r = 0$  and  $r = \infty$
- FIND GREEN FUNCTION RESULT

# Conclusion

- First steps in understanding spatially anisotropic AdS/CFT
- Several obvious next steps
  - Flows between metrics
  - Full numerical computation
  - How do we understand unusual behavior near origin?

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